

Conduction through Multilayers Plane Wall

Introduction:

In heat transfer analysis, we are often interested in the rate of *heat transfer through a medium under steady conditions* and surface temperatures. Such problems can be solved easily without involving any differential equations by the introduction of *thermal resistance concepts* in an analogous manner to *electrical circuit* problems.

In this case, the *thermal resistance corresponds to electrical resistance, temperature difference corresponds to voltage, and the heat transfer rate corresponds to electric current*. We start with *one-dimensional steady heat conduction* in a *multilayer plane wall, cylinder, and sphere*, and develop relations for *thermal resistances* in these geometries.

Consider a *plane wall of thickness L* and average thermal conductivity *k*. The two surfaces of the wall are maintained at constant temperatures of T_1 and T_2 . For one-dimensional steady heat conduction through the wall, we have $T(x)$. Then *Fourier's law* of heat conduction for the wall can be expressed as:

$$\dot{Q}_{\text{cond, wall}} = -kA \frac{dT}{dx} \quad (W)$$

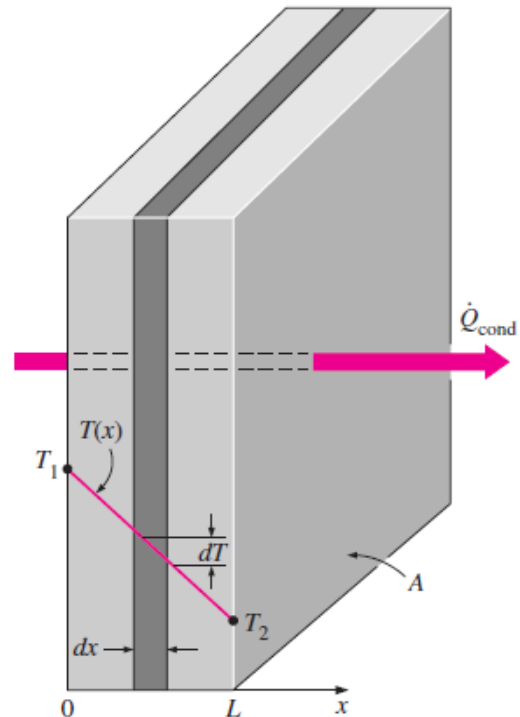
where the rate of conduction heat transfer $\dot{Q}_{\text{cond, wall}}$ and the wall area A are constant. Thus we have $dT/dx = \text{constant}$, which means that *the temperature through the wall varies linearly with x*. That is, the temperature distribution in the wall under steady conditions is a *straight line*. Separating the variables in the above equation and integrating from $x = 0$, where $T(0) = T_1$, to $x = L$, where $T(L) = T_2$, we get:

$$\int_{x=0}^L \dot{Q}_{\text{cond, wall}} dx = - \int_{T=T_1}^{T_2} kA dT$$

$$\dot{Q}_{\text{cond, wall}} [x]_0^L = -kA. [T]_{T_1}^{T_2}$$

Performing the *integrations* and rearranging gives

$$\dot{Q}_{\text{cond, wall}} = kA \frac{T_1 - T_2}{L} \quad (W)$$



The Thermal Resistance Concept

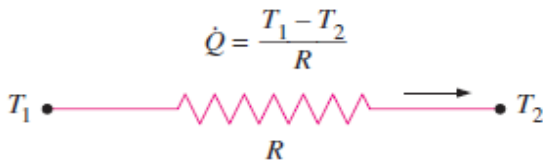
Heat conduction through a plane wall can be rearranged as

$$\dot{Q}_{cond,wall} = \frac{T_1 - T_2}{R_{wall}} \quad (W)$$

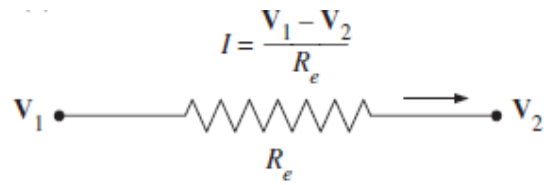
Where: $R_{wall} = \frac{L}{kA}$ ($^{\circ}C/W$) is the **thermal resistance** of the wall against heat conduction or simply the **conduction resistance** of the wall.

The equation above for heat flow is analogous to the relation for **electric current flow I**, expressed as]

$$I = \frac{V_1 - V_2}{R_e}$$



(a) Heat flow



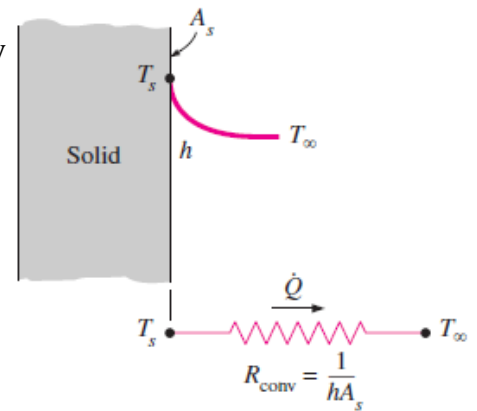
(b) Electric current flow

Consider convection heat transfer from a solid surface of area A_s and temperature T_s to a fluid whose temperature sufficiently far from the surface is T_{∞} , with a convection heat transfer coefficient h .

Newton's law of cooling for convection heat transfer rate:

$\dot{Q}_{conv} = hA_s(T_s - T_{\infty})$ can be rearranged as:

$$\dot{Q}_{conv} = \frac{T_s - T_{\infty}}{R_{conv}} \quad (W) \quad \text{where} \quad R_{conv} = \frac{1}{hA_s} \quad (^{\circ}C/W)$$

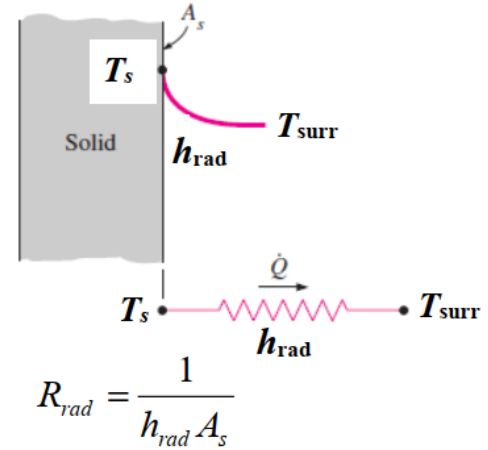


is the **thermal resistance** of the surface against heat convection, or simply the **convection resistance** of the surface.

The **rate of radiation heat transfer** between a surface of emissivity ϵ and area A_s at temperature T_s and the surrounding surfaces at some average temperature T_{surr} can be expressed as:

$$\dot{Q}_{rad} = \epsilon \sigma A_s (T_s^4 - T_{surr}^4) = \epsilon \sigma A_s (T_s^2 - T_{surr}^2)(T_s^2 + T_{surr}^2)$$

$$\begin{aligned}\dot{Q}_{rad} &= \varepsilon\sigma A_s (T_s - T_{surr})(T_s + T_{surr})(T_s^2 + T_{surr}^2) \\ &= A_s (T_s - T_{surr}) \underbrace{\varepsilon\sigma (T_s + T_{surr})(T_s^2 + T_{surr}^2)}_{h_{rad}} \\ \dot{Q}_{rad} &= h_{rad} A_s (T_s - T_{surr}) = \frac{T_s - T_{surr}}{R_{rad}} \quad (W)\end{aligned}$$



Where: $h_{rad} = \varepsilon\sigma (T_s + T_{surr})(T_s^2 + T_{surr}^2) \quad (W/m^2 \cdot K)$

$R_{rad} = \frac{1}{h_{rad} A_s} \quad (K/W)$ is the **thermal resistance** of a surface against radiation, or the

radiation resistance, and h_{rad} is the **radiation heat transfer coefficient**.

Emissivity (ε): 0 – 1, **Stefan Boltzmann constant: $5.67 \times 10^{-8} \quad (W/m^2 \cdot K)$**

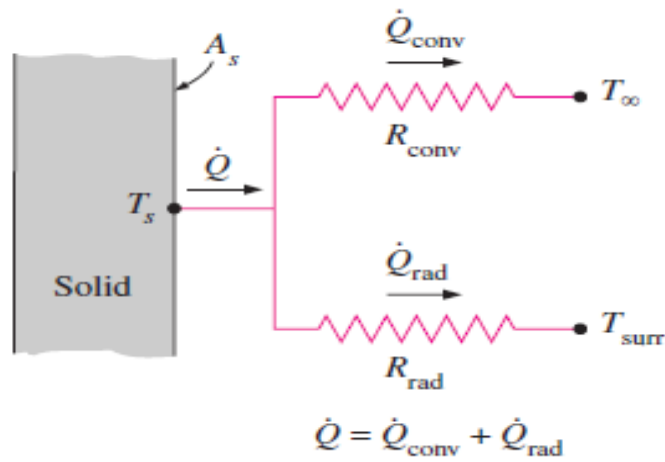
Note that both T_s and T_{surr} *must be in K* in the evaluation of h_{rad} . The definition of the radiation heat transfer coefficient enables us to express radiation conveniently in an analogous manner to convection in terms of a temperature difference. **But h_{rad} depends strongly on temperature while h_{conv} usually does not.**

A surface exposed to the surrounding air involves **convection** and **radiation** simultaneously, The convection and radiation resistances are **parallel** to each other, as shown in the Fig., and may cause some complication in the thermal resistance network.

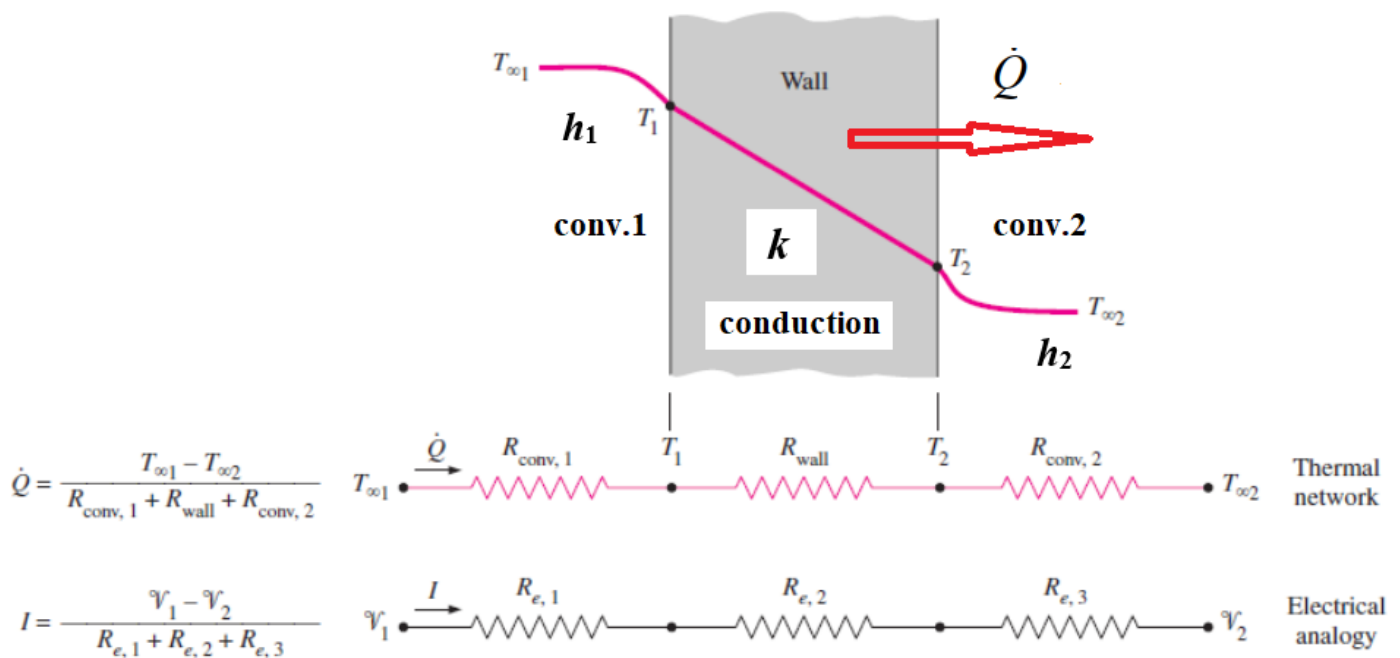
When $T_{surr} = T_\infty$, the radiation effect can properly be accounted for by replacing h in the convection resistance relation by:

$$h_{combined} = h_{conv} + h_{rad} \quad (W/m^2 \cdot K)$$

where $h_{combined}$ is the **combined heat transfer coefficient**. In this way, all the complications associated with radiation are avoided.



Thermal Resistance Network:



The thermal resistance network for heat transfer through a plane wall subjected to convection on both sides, and the electrical analogy.

Now consider steady one-dimensional heat flow through a plane wall of thickness L , area A , and thermal conductivity k that is exposed to convection on both sides to fluids at temperatures $T_{\infty 1}$ and $T_{\infty 2}$ with heat transfer coefficients h_1 and h_2 , respectively, as shown in the Figure. Assuming $T_{\infty 2} < T_{\infty 1}$, the variation of temperature will be as shown in the figure. Note that the temperature varies linearly in the wall, and asymptotically approaches $T_{\infty 1}$ and $T_{\infty 2}$ in the fluids as we move away from the wall.

Under steady conditions we have:

or
$$\left(\begin{array}{c} \text{Rate of} \\ \text{heat convection} \\ \text{into the wall} \end{array} \right) = \left(\begin{array}{c} \text{Rate of} \\ \text{heat conduction} \\ \text{through the wall} \end{array} \right) = \left(\begin{array}{c} \text{Rate of} \\ \text{heat convection} \\ \text{from the wall} \end{array} \right)$$

$$\dot{Q} = h_1 A(T_{\infty 1} - T_1) = kA \frac{T_1 - T_2}{L} = h_2 A(T_2 - T_{\infty 2})$$

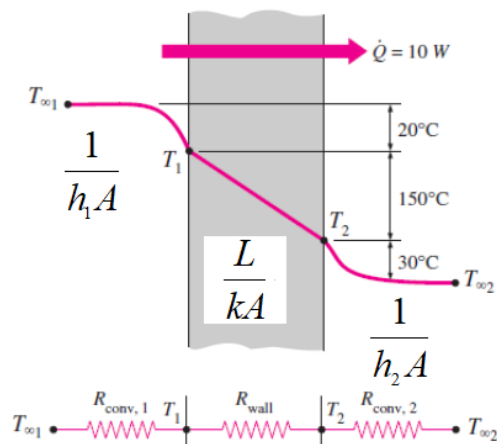
$$\dot{Q} = \frac{T_{\infty 1} - T_1}{1/h_1 A} = \frac{T_1 - T_2}{L/kA} = \frac{T_2 - T_{\infty 2}}{1/h_2 A} = \frac{T_{\infty 1} - T_1}{R_{\text{conv}, 1}} = \frac{T_1 - T_2}{R_{\text{wall}}} = \frac{T_2 - T_{\infty 2}}{R_{\text{conv}, 2}}$$

Adding the numerators and denominators yields,

$$\dot{Q} = \frac{T_{\infty 1} - T_{\infty 2}}{R_{\text{total}}} \quad (\text{W})$$

Where
$$R_{\text{total}} = R_{\text{conv}, 1} + R_{\text{wall}} + R_{\text{conv}, 2} = \frac{1}{h_1 A} + \frac{L}{kA} + \frac{1}{h_2 A} \quad (^\circ\text{C}/\text{W})$$

$$\dot{Q} = \frac{\Delta T}{R_{\text{total}}}$$



Overall heat transfer coefficient (U)

It is sometimes *heat transfer through a medium* convenient to express in an analogous manner to Newton's law of cooling as,

$$\dot{Q} = UA \Delta T \quad (\text{W})$$

where U is the *overall heat transfer coefficient* ($\text{W}/\text{m}^2 \cdot \text{K}$).

$$UA = \frac{1}{R_{\text{total}}} \quad (\text{W}), \text{ and } R_{\text{total}} = \frac{1}{UA}$$

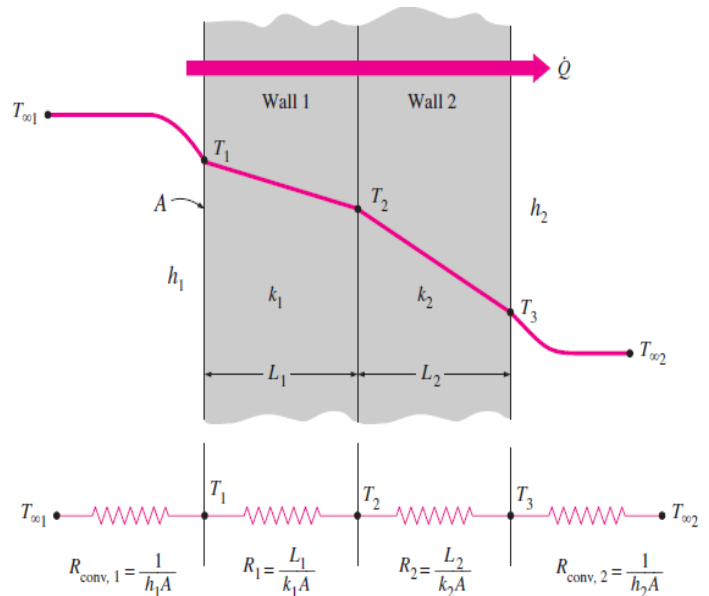
Multilayer Plane Walls

In practice we often encounter plane walls that consist of *several layers of different materials*. The thermal resistance concept can still be used to determine the rate of steady heat transfer through such *composite walls*.

$$\dot{Q} = \frac{T_{\infty 1} - T_{\infty 2}}{R_{total}} \quad (W)$$

R_{total} = the **total thermal resistance**, expressed as:

$$\begin{aligned} R_{total} &= R_{conv,1} + R_{wall,1} + R_{wall,2} + R_{conv,2} \\ &= \frac{1}{h_1 A} + \frac{L_1}{k_1 A} + \frac{L_2}{k_2 A} + \frac{1}{h_2 A} \quad (^\circ C/W) \end{aligned}$$



Once \dot{Q} is **known**, an unknown surface temperature T_j at any surface or interface j can be determined from:

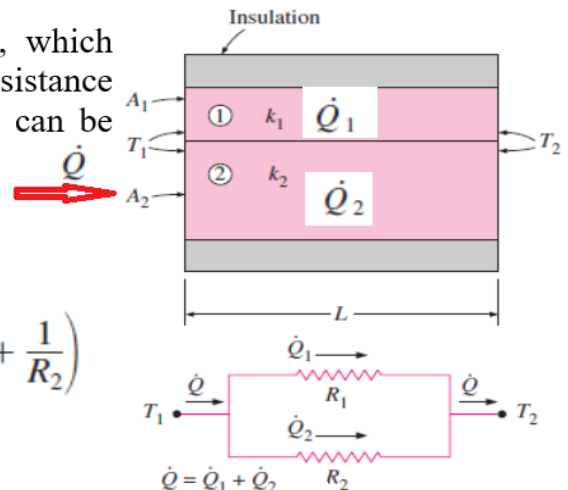
$$\dot{Q} = \frac{T_i - T_j}{R_{total, i-j}} \quad \text{where } T_i \text{ is a known temperature at location } i \text{ and } R_{total, i-j} \text{ is the total thermal resistance between locations } i \text{ and } j.$$

$$\text{To find } T_1: \dot{Q} = \frac{T_{\infty 1} - T_1}{R_{conv,1}} \quad \text{To find } T_2: \dot{Q} = \frac{T_{\infty 1} - T_2}{R_{conv,1} + R_1} \quad \text{To find } T_3: \dot{Q} = \frac{T_3 - T_{\infty 2}}{R_{conv,2}}$$

Generalized Thermal Resistance Networks

The **thermal resistance** concept or the **electrical analogy** can also be used to solve steady heat transfer problems that involve parallel layers or **combined series-parallel arrangements**.

Consider the **composite wall** shown in the Figure, which consists of **two parallel layers**. The thermal resistance network, which consists of two parallel resistances, can be represented as shown in the figure.



$$\dot{Q} = \dot{Q}_1 + \dot{Q}_2 = \frac{T_1 - T_2}{R_1} + \frac{T_1 - T_2}{R_2} = (T_1 - T_2) \left(\frac{1}{R_1} + \frac{1}{R_2} \right)$$

Utilizing **electrical analogy**, we get:

$$\dot{Q} = \frac{T_1 - T_2}{R_{total}} \quad \text{where} \quad \frac{1}{R_{total}} = \frac{1}{R_1} + \frac{1}{R_2} = \frac{R_1 + R_2}{R_1 R_2}$$

$$\rightarrow R_{total} = \frac{R_1 R_2}{R_1 + R_2}$$

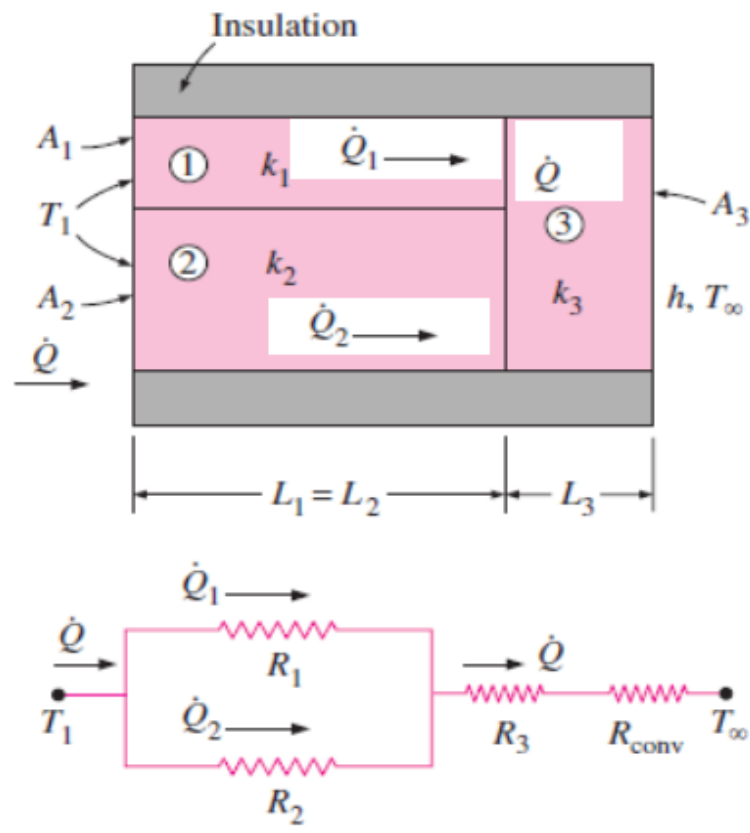
Now consider the **combined series-parallel arrangement** shown, the total rate of heat transfer through this composite system can again be expressed as:

$$\dot{Q} = \frac{T_1 - T_\infty}{R_{total}}$$

$$R_{total} = R_{12} + R_3 + R_{conv} = \frac{R_1 R_2}{R_1 + R_2} + R_3 + R_{conv}$$

$$R_1 = \frac{L_1}{k_1 A_1}, \quad R_2 = \frac{L_2}{k_2 A_2}, \quad R_3 = \frac{L_3}{k_3 A_3}, \quad R_{conv} = \frac{1}{h A_3}$$

Once the **individual thermal resistances** are evaluated, the **total resistance and the total rate of heat transfer** can easily be determined from the relations above.



EXAMPLE: Heat Loss through Double-Pane Windows

Consider a 0.8 m high and 1.5 m wide double-pane window consisting of two 4 mm thick layers of glass ($k= 0.78 \text{ W/m}\cdot^\circ\text{C}$) separated by a 10 mm wide stagnant air space ($k= 0.026 \text{ W/m}\cdot^\circ\text{C}$). **Determine the steady rate of heat transfer through this double-pane window and the temperature of its inner surface** for a day during which the room is maintained at 20°C while the temperature of the outdoors is -10°C . Take the convection heat transfer coefficients on the inner and outer surfaces of the window to be $h_1= 10 \text{ W/m}^2 \cdot ^\circ\text{C}$ and $h_2= 40 \text{ W/m}^2\cdot^\circ\text{C}$, which includes the effects of radiation.

Solution:

Area: $A = 0.8 \times 1.5 = 1.2 \text{ m}^2$

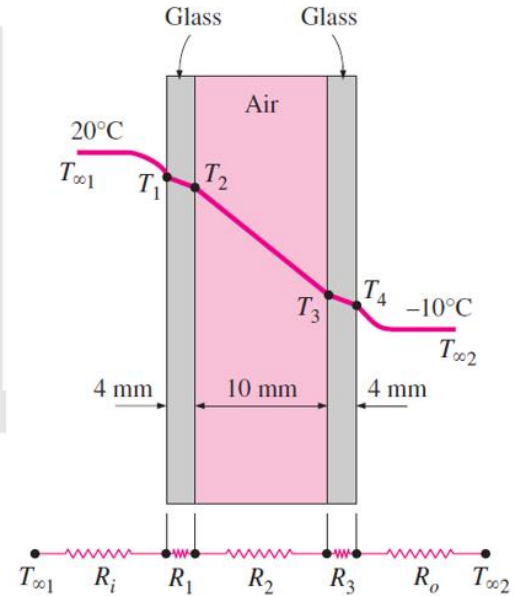
$$R_i = R_{\text{conv},1} = \frac{1}{h_1 A} = \frac{1}{(10 \text{ W/m}^2 \cdot ^\circ\text{C})(1.2 \text{ m}^2)} = 0.08333^\circ\text{C/W}$$

$$R_1 = R_3 = R_{\text{glass}} = \frac{L_1}{k_1 A} = \frac{0.004 \text{ m}}{(0.78 \text{ W/m} \cdot ^\circ\text{C})(1.2 \text{ m}^2)} = 0.00427^\circ\text{C/W}$$

$$R_2 = R_{\text{air}} = \frac{L_2}{k_2 A} = \frac{0.01 \text{ m}}{(0.026 \text{ W/m} \cdot ^\circ\text{C})(1.2 \text{ m}^2)} = 0.3205^\circ\text{C/W}$$

$$R_o = R_{\text{conv},2} = \frac{1}{h_2 A} = \frac{1}{(40 \text{ W/m}^2 \cdot ^\circ\text{C})(1.2 \text{ m}^2)} = 0.02083^\circ\text{C/W}$$

Noting that all three resistances are in series, the total resistance is



$$\begin{aligned} R_{\text{total}} &= R_{\text{conv},1} + R_{\text{glass},1} + R_{\text{air}} + R_{\text{glass},2} + R_{\text{conv},2} \\ &= 0.08333 + 0.00427 + 0.3205 + 0.00427 + 0.02083 \\ &= 0.4332^\circ\text{C/W} \end{aligned}$$

Then the steady rate of heat transfer through the window becomes

$$\dot{Q} = \frac{T_{\infty 1} - T_{\infty 2}}{R_{\text{total}}} = \frac{[20 - (-10)]^\circ\text{C}}{0.4332^\circ\text{C/W}} = \mathbf{69.2 \text{ W}}$$

The inner surface temperature of the window in this case will be

$$\dot{Q} = \frac{T_{\infty 1} - T_1}{R_{\text{conv},1}}$$

$$T_1 = T_{\infty 1} - \dot{Q} R_{\text{conv},1} = 20^\circ\text{C} - (69.2 \text{ W})(0.08333^\circ\text{C/W}) = \mathbf{14.2^\circ\text{C}}$$